

Proceedings of the 12th International Conference on
Computational Fluid Dynamics in the Oil & Gas,
Metallurgical and Process Industries

Progress in Applied CFD – CFD2017



SINTEF Proceedings

Editors:

Jan Erik Olsen and Stein Tore Johansen

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PREFACE

This book contains all manuscripts approved by the reviewers and the organizing committee of the 12th International Conference on Computational Fluid Dynamics in the Oil & Gas, Metallurgical and Process Industries. The conference was hosted by SINTEF in Trondheim in May/June 2017 and is also known as CFD2017 for short. The conference series was initiated by CSIRO and Phil Schwarz in 1997. So far the conference has been alternating between CSIRO in Melbourne and SINTEF in Trondheim. The conferences focuses on the application of CFD in the oil and gas industries, metal production, mineral processing, power generation, chemicals and other process industries. In addition pragmatic modelling concepts and bio-mechanical applications have become an important part of the conference. The papers in this book demonstrate the current progress in applied CFD.

The conference papers undergo a review process involving two experts. Only papers accepted by the reviewers are included in the proceedings. 108 contributions were presented at the conference together with six keynote presentations. A majority of these contributions are presented by their manuscript in this collection (a few were granted to present without an accompanying manuscript).

The organizing committee would like to thank everyone who has helped with review of manuscripts, all those who helped to promote the conference and all authors who have submitted scientific contributions. We are also grateful for the support from the conference sponsors: ANSYS, SFI Metal Production and NanoSim.

Stein Tore Johansen & Jan Erik Olsen



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CONTENTS

PRAGMATIC MODELLING	9
On pragmatism in industrial modeling. Part III: Application to operational drilling	11
CFD modeling of dynamic emulsion stability	23
Modelling of interaction between turbines and terrain wakes using pragmatic approach	29
FLUIDIZED BED	37
Simulation of chemical looping combustion process in a double looping fluidized bed reactor with cu-based oxygen carriers.....	39
Extremely fast simulations of heat transfer in fluidized beds.....	47
Mass transfer phenomena in fluidized beds with horizontally immersed membranes	53
A Two-Fluid model study of hydrogen production via water gas shift in fluidized bed membrane reactors	63
Effect of lift force on dense gas-fluidized beds of non-spherical particles	71
Experimental and numerical investigation of a bubbling dense gas-solid fluidized bed	81
Direct numerical simulation of the effective drag in gas-liquid-solid systems	89
A Lagrangian-Eulerian hybrid model for the simulation of direct reduction of iron ore in fluidized beds.....	97
High temperature fluidization - influence of inter-particle forces on fluidization behavior	107
Verification of filtered two fluid models for reactive gas-solid flows	115
BIOMECHANICS.....	123
A computational framework involving CFD and data mining tools for analyzing disease in carotid artery	125
Investigating the numerical parameter space for a stenosed patient-specific internal carotid artery model.....	133
Velocity profiles in a 2D model of the left ventricular outflow tract, pathological case study using PIV and CFD modeling.....	139
Oscillatory flow and mass transport in a coronary artery.....	147
Patient specific numerical simulation of flow in the human upper airways for assessing the effect of nasal surgery.....	153
CFD simulations of turbulent flow in the human upper airways	163
OIL & GAS APPLICATIONS	169
Estimation of flow rates and parameters in two-phase stratified and slug flow by an ensemble Kalman filter	171
Direct numerical simulation of proppant transport in a narrow channel for hydraulic fracturing application	179
Multiphase direct numerical simulations (DNS) of oil-water flows through homogeneous porous rocks	185
CFD erosion modelling of blind tees	191
Shape factors inclusion in a one-dimensional, transient two-fluid model for stratified and slug flow simulations in pipes	201
Gas-liquid two-phase flow behavior in terrain-inclined pipelines for wet natural gas transportation	207

NUMERICS, METHODS & CODE DEVELOPMENT	213
Innovative computing for industrially-relevant multiphase flows	215
Development of GPU parallel multiphase flow solver for turbulent slurry flows in cyclone.....	223
Immersed boundary method for the compressible Navier–Stokes equations using high order summation-by-parts difference operators	233
Direct numerical simulation of coupled heat and mass transfer in fluid-solid systems	243
A simulation concept for generic simulation of multi-material flow, using staggered Cartesian grids.....	253
A cartesian cut-cell method, based on formal volume averaging of mass, momentum equations.....	265
SOFT: a framework for semantic interoperability of scientific software	273
POPULATION BALANCE	279
Combined multifluid-population balance method for polydisperse multiphase flows	281
A multifluid-PBE model for a slurry bubble column with bubble size dependent velocity, weight fractions and temperature.....	285
CFD simulation of the droplet size distribution of liquid-liquid emulsions in stirred tank reactors	295
Towards a CFD model for boiling flows: validation of QMOM predictions with TOPFLOW experiments	301
Numerical simulations of turbulent liquid-liquid dispersions with quadrature-based moment methods.....	309
Simulation of dispersion of immiscible fluids in a turbulent couette flow	317
Simulation of gas-liquid flows in separators - a Lagrangian approach.....	325
CFD modelling to predict mass transfer in pulsed sieve plate extraction columns	335
BREAKUP & COALESCENCE	343
Experimental and numerical study on single droplet breakage in turbulent flow	345
Improved collision modelling for liquid metal droplets in a copper slag cleaning process	355
Modelling of bubble dynamics in slag during its hot stage engineering.....	365
Controlled coalescence with local front reconstruction method	373
BUBBLY FLOWS	381
Modelling of fluid dynamics, mass transfer and chemical reaction in bubbly flows	383
Stochastic DSMC model for large scale dense bubbly flows.....	391
On the surfacing mechanism of bubble plumes from subsea gas release.....	399
Bubble generated turbulence in two fluid simulation of bubbly flow	405
HEAT TRANSFER	413
CFD-simulation of boiling in a heated pipe including flow pattern transitions using a multi-field concept	415
The pear-shaped fate of an ice melting front	423
Flow dynamics studies for flexible operation of continuous casters (flow flex cc).....	431
An Euler-Euler model for gas-liquid flows in a coil wound heat exchanger.....	441
NON-NEWTONIAN FLOWS.....	449
Viscoelastic flow simulations in disordered porous media	451
Tire rubber extrudate swell simulation and verification with experiments	459
Front-tracking simulations of bubbles rising in non-Newtonian fluids.....	469
A 2D sediment bed morphodynamics model for turbulent, non-Newtonian, particle-loaded flows.....	479

METALLURGICAL APPLICATIONS.....	491
Experimental modelling of metallurgical processes	493
State of the art: macroscopic modelling approaches for the description of multiphysics phenomena within the electroslag remelting process	499
LES-VOF simulation of turbulent interfacial flow in the continuous casting mold	507
CFD-DEM modelling of blast furnace tapping	515
Multiphase flow modelling of furnace tapholes	521
Numerical predictions of the shape and size of the raceway zone in a blast furnace.....	531
Modelling and measurements in the aluminium industry - Where are the obstacles?	541
Modelling of chemical reactions in metallurgical processes.....	549
Using CFD analysis to optimise top submerged lance furnace geometries	555
Numerical analysis of the temperature distribution in a martensitic stainless steel strip during hardening.....	565
Validation of a rapid slag viscosity measurement by CFD.....	575
Solidification modeling with user defined function in ANSYS Fluent.....	583
Cleaning of polycyclic aromatic hydrocarbons (PAH) obtained from ferroalloys plant.....	587
Granular flow described by fictitious fluids: a suitable methodology for process simulations	593
A multiscale numerical approach of the dripping slag in the coke bed zone of a pilot scale Si-Mn furnace.....	599
INDUSTRIAL APPLICATIONS	605
Use of CFD as a design tool for a phosphoric acid plant cooling pond	607
Numerical evaluation of co-firing solid recovered fuel with petroleum coke in a cement rotary kiln: Influence of fuel moisture	613
Experimental and CFD investigation of fractal distributor on a novel plate and frame ion-exchanger	621
COMBUSTION	631
CFD modeling of a commercial-size circle-draft biomass gasifier.....	633
Numerical study of coal particle gasification up to Reynolds numbers of 1000.....	641
Modelling combustion of pulverized coal and alternative carbon materials in the blast furnace raceway	647
Combustion chamber scaling for energy recovery from furnace process gas: waste to value	657
PACKED BED.....	665
Comparison of particle-resolved direct numerical simulation and 1D modelling of catalytic reactions in a packed bed	667
Numerical investigation of particle types influence on packed bed adsorber behaviour	675
CFD based study of dense medium drum separation processes	683
A multi-domain 1D particle-reactor model for packed bed reactor applications.....	689
SPECIES TRANSPORT & INTERFACES	699
Modelling and numerical simulation of surface active species transport - reaction in welding processes	701
Multiscale approach to fully resolved boundary layers using adaptive grids.....	709
Implementation, demonstration and validation of a user-defined wall function for direct precipitation fouling in Ansys Fluent.....	717

FREE SURFACE FLOW & WAVES	727
Unresolved CFD-DEM in environmental engineering: submarine slope stability and other applications.....	729
Influence of the upstream cylinder and wave breaking point on the breaking wave forces on the downstream cylinder	735
Recent developments for the computation of the necessary submergence of pump intakes with free surfaces	743
Parallel multiphase flow software for solving the Navier-Stokes equations	752
 PARTICLE METHODS	 759
A numerical approach to model aggregate restructuring in shear flow using DEM in Lattice-Boltzmann simulations	761
Adaptive coarse-graining for large-scale DEM simulations.....	773
Novel efficient hybrid-DEM collision integration scheme.....	779
Implementing the kinetic theory of granular flows into the Lagrangian dense discrete phase model.....	785
Importance of the different fluid forces on particle dispersion in fluid phase resonance mixers	791
Large scale modelling of bubble formation and growth in a supersaturated liquid.....	798
 FUNDAMENTAL FLUID DYNAMICS	 807
Flow past a yawed cylinder of finite length using a fictitious domain method	809
A numerical evaluation of the effect of the electro-magnetic force on bubble flow in aluminium smelting process.....	819
A DNS study of droplet spreading and penetration on a porous medium.....	825
From linear to nonlinear: Transient growth in confined magnetohydrodynamic flows.....	831

MULTIPHASE DIRECT NUMERICAL SIMULATIONS (DNS) OF OIL-WATER FLOWS THROUGH HOMOGENEOUS POROUS ROCKS

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ABSTRACT

Water flooding is commonly used to recover oil from porous rocks using pressurized water. Present study focuses on fully resolved pore-scale level multiphase Direct Numerical Simulations (DNS) of oil-water flows through homogeneous porous rocks. A second order accurate implicit Immersed Boundary Method (IBM) is used to resolve fluid-solid interactions on a non-body fitted Cartesian computational grid. Dynamic evolution of the fluid-fluid interface is tracked by a mass conservative sharp interface Volume of Fluid (VOF) method. The IBM and VOF method are coupled by a prescribed contact angle boundary condition at the fluid-fluid-solid contact line. Our method has been extensively validated especially for the test cases involving oil-water flows. Simulations of water flooding process through periodic homogeneous configurations of spheres are performed based on typical pore-scale capillary and Reynolds numbers. Effect of wettability on the mobility of oil through oil-wet and neutrally-wet rocks has been quantified as well.

Keywords: water flooding, porous rocks, Immersed Boundary Method (IBM), Volume of Fluid (VOF), contact angle .

NOMENCLATURE

Notations

p	Pressure, [Pa]
F	Fluid phase fraction, [-]
M	Dynamic viscosity ratio, [-]
Re	Reynolds number, [-]
Ca	Capillary number, [-]
S	Fluid phase saturation, [-]
k	Permeability, [m^2]
x	x co-ordinate, [m]
y	y co-ordinate, [m]
z	z co-ordinate, [m]

Greek Symbols

ρ	Mass density, [kg/m^3]
μ	Dynamic viscosity, [$Pa \cdot s$]
σ	Surface tension, [N/m]
Δ	Grid size, [m]
θ	(Static) contact angle, [$degree$]
ϕ	Porosity, [-]

Vectors

\mathbf{u}	Velocity, [m/s]
\mathbf{F}_σ	Surface tension force, [N]
\mathbf{g}	Gravitational acceleration, [m/s^2]

Sub/superscripts

1,2	Fluid phase number
w	Wetting fluid
nw	Non-wetting fluid

INTRODUCTION

Multiphase flows in complex structures are encountered widely in nature and technology. One such example is water flooding (Sheng, 2014) used for oil recovery. After the primary (natural) recovery of the oil from reservoir, large amount of oil remains trapped in the porous rocks. Secondary and ternary recovery processes (e.g. water flooding, gas injection, thermal processes, chemical flooding etc.) are then used for further recovery of such residual oil. In water flooding high pressure and/or high temperature water is pushed through porous rocks which carries oil out of the porous bed. Focus of the current work is to perform pore-scale simulations of oil-water multiphase flows through complex rock structures during such a water flooding process. To serve this purpose, three different problems need to be tackled: i) oil-water multi-fluid interface tracking, ii) interactions between fluids (oil or water) and complex solid geometries, and iii) three-phase contact line dynamics.

A wide range of numerical methods (e.g. Front Tracking, Volume of Fluid, Level Set etc.) has been developed and tested successfully to track multi-fluid interfaces (Wörner, 2003). They all differ with respect to tackling the following complexity: i) interface advection based on local velocity field and ii) surface tension force based on local interface curvature. The volume of fluid (VOF) method (Hirt and Nichols, 1981; Youngs, 1982) uses a color function F which denotes the local fluid phase fraction in the immiscible mixture of fluids. Advection of F is governed by pseudo-Lagrangian geometrical advection schemes to minimize numerical diffusion. This particular feature makes VOF the most mass conservative among all multi-fluid interface tracking methods. The density-scaled continuum surface force (CSF) model proposed by Brackbill *et al.* (1992) is used to evaluate the surface tension force for its simplicity and robustness especially in the presence of complex solid boundaries.

The immersed boundary method (IBM) (Mittal and Iaccarino, 2005) is a set of computational techniques which

uses non-body fitted (mainly Cartesian) grids for simulating fluid-solid interactions through complex geometries. It eliminates the need of traditional unstructured body fitted grids and hence provides simplicity in grid generation and discretization, ease of code development, less memory requirement and higher computational efficiency. A direct forcing approach used in IBM produces a sharp fluid-solid interface without spatial spreading. So, in the present work a direct forcing, implicit, second order IBM (Deen *et al.*, 2012; Das *et al.*, 2016) is used which does not require any calibration for the different complex geometries.

Fluid-fluid interfaces in contact with solid boundaries produce a three-phase contact line. The contact line behavior is determined by microscopic physicochemical interactions between molecules of two different immiscible fluids and the solid substrate, and it can drastically affect the statics and dynamics of the bulk flow (Snoeijer and Andreotti, 2013). An apparent contact angle at the macroscopic length scales is a result of these microscopic interactions at the contact line. In the present work, a coupled IBM-VOF method has been used to simulate contact line dynamics.

This paper deals with validation/verification and application of the developed IBM-VOF solver specifically for multiphase flows involving oil and water. It is organized as follows: We first describe the governing Navier-Stokes equations for multiphase flows along with the F -advection equation. Next, we discuss numerical and implementation details in brief. Further, the coupled IBM-VOF solver is extensively validated/verified for the test cases involving oil-water flows. Last, a water flooding process is simulated on the pore-scale and the effect of wettability on the mobility of oil has been quantified for oil-wet and neutrally-wet homogeneous rocks.

MODEL DESCRIPTION

Governing equations

For incompressible multiphase flows the Navier-Stokes equations can be combined into a single equation for \mathbf{u} in the entire domain. Surface tension due to the presence of a curved deformable fluid-fluid interface is taken into account by a local volumetric \mathbf{F}_σ . The governing mass and momentum conservation equations for unsteady, incompressible, Newtonian, multiphase flows are expressed as follows:

$$\nabla \cdot \mathbf{u} = 0 \quad (1)$$

$$\rho \frac{\partial \mathbf{u}}{\partial t} + \rho \nabla \cdot (\mathbf{u}\mathbf{u}) = -\nabla p + \nabla \cdot \boldsymbol{\tau} + \rho \mathbf{g} + \mathbf{F}_\sigma \quad (2)$$

where $\boldsymbol{\tau} = \mu[\nabla \mathbf{u} + (\nabla \mathbf{u})^T]$ is the fluid stress tensor. Advection of F is governed by the following equation:

$$\frac{DF}{Dt} = \frac{\partial F}{\partial t} + \mathbf{u} \cdot \nabla F = 0 \quad (3)$$

This equation expresses that the interface is advected with the local fluid velocity. To evaluate the local averaged density, linear averaging of the densities of the fluid 1 ($F = 1$) and fluid 2 ($F = 0$) is used:

$$\rho = F\rho_1 + (1-F)\rho_2 \quad (4)$$

Similarly, the local average μ could also be evaluated by linear averaging of the dynamic viscosities of individual fluid phase. Alternatively, following a fundamental approach proposed by Prosperetti (2002), the local average kinematic viscosity is evaluated by harmonic averaging of the kinematic viscosities of the individual fluid phases:

$$\frac{\rho}{\mu} = F \frac{\rho_1}{\mu_1} + (1-F) \frac{\rho_2}{\mu_2} \quad (5)$$

In all the computations reported in this paper, Eq. (4) and Eq. (5) are used to evaluate the local average density and dynamic viscosity, respectively.

Numerical and Implementation Details

In the current implementation, the finite volume method is used to solve the mass and momentum conservation equations in each control volume of a staggered computational grid. To simulate multiphase flows involving multi fluid-solid interactions three things need to be dealt with: i) dynamic interactions between fluids and non-deformable solids, ii) dynamics of deformable fluid-fluid interfaces, and iii) contact line dynamics at the three phase contact lines.

An implicit (direct) second-order accurate IBM proposed by Deen *et al.* (2012) and described in detail by Das *et al.* (2016) has been used to apply no-slip boundary conditions for fluids-solid interactions at immersed boundaries. VOF (Van Sint Annaland *et al.*, 2005) to track fluid-fluid interface consists of three main parts: i) solution of F -advection equation, ii) computation of the volumetric \mathbf{F}_σ , and iii) smoothing of F . In presence of fluid-fluid interfaces with immersed solids, contact line dynamics plays a major role in wetting-dewetting phenomena. It is incorporated in our coupled IBM-VOF framework (Patel *et al.*, 2017) by applying the apparent contact angle as a boundary condition at contact lines. The contact angle may have a single value (static) or different values (dynamic) depending upon the local contact parameters and fluid properties. The effect of the contact angle is taken into account by modifying the interface normals at the solid boundaries. Readers may refer to the mentioned literature for intrinsic numerical and implementation details.

VERIFICATION AND VALIDATION

The IBM and VOF implementations have been tested individually by Das *et al.* (2016) and Van Sint Annaland *et al.* (2005) respectively for various cases and they obtained excellent agreement with numerical and experimental results published in literature. The 3D coupled IBM-VOF implementation with contact line dynamics has been validated by Patel *et al.* (2017) with static and/or dynamic contact angles for i) the equilibrium shape of a droplet on a flat surface, ii) the equilibrium shape of a droplet on a spherical surface, and iii) temporal evolution of droplet contact radius on flat surface (generated with and without IBM).

In this section, we investigate additional validation/verification test cases specifically addressing the oil-water multiphase flows. First, a test case of 2D multiphase Poiseuille flow is considered to compare the velocity profiles and relative permeabilities with analytical expressions. Further, a viscous fingering phenomenon is simulated in a 2D channel and finger characteristics parameters i.e. dimensionless finger width and tip radius have been compared with results published in literature. Last, a grid independence study for the contact force is presented for 3D coupled IBM-VOF implementation.

2D Multiphase Poiseuille Flow

In this section, we investigate the co-current multiphase Poiseuille flow of two immiscible fluids namely wetting and non-wetting fluid. The wetting fluid is in contact with the channel walls ($a < |y| < L$) whereas the non-wetting fluid resides between the layers of wetting fluid ($0 < |y| < a$) as shown in Figure 1. No-slip boundary conditions are applied at the

channel walls ($|y| = L$) and the flow in x -direction is periodic. Both fluids have the same density ($\rho_w = \rho_{nw}$) and different dynamic viscosities giving the definition of $M = \mu_{nw}/\mu_w$. A constant body force G is applied to both fluids in the x -direction such as the flow remains in the Stokes regime ($Re \ll 1$). Due to the existence of a flat fluid-fluid interface, the surface tension force doesn't play any role in this problem.

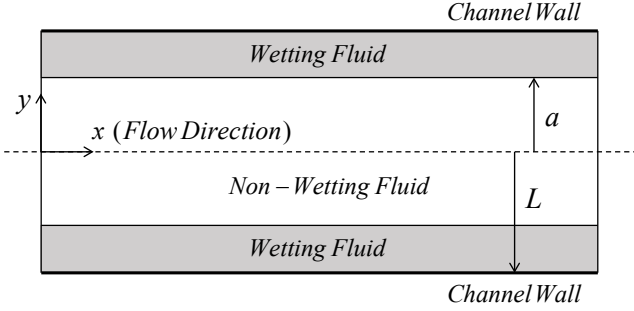


Figure 1: Schematic diagram of 2D multiphase Poiseuille flow.

For the given value of wetting and non-wetting saturation ($S_w = (1 - a)/L$ and $S_{nw} = a/L$), the analytical expressions for the velocity of the wetting and non-wetting fluids (u_w and u_{nw}) are given by:

$$\begin{aligned} u_w(y) &= \frac{G}{2\mu_w}(L^2 - y^2) \\ u_{nw}(y) &= \frac{G}{2\mu_w}(L^2 - a^2) + \frac{G}{2\mu_{nw}}(a^2 - y^2) \end{aligned} \quad (6)$$

Similarly, the analytical expressions for the relative permeability of each fluid ($k_{r,w}$ and $k_{r,nw}$) can be given as,

$$\begin{aligned} k_{r,w} &= \frac{1}{2}S_w^2(3 - S_w) \\ k_{r,nw} &= S_{nw} \left[\frac{3}{2}M + S_{nw}^2 \left(1 - \frac{3}{2}M \right) \right] \end{aligned} \quad (7)$$

Readers are referred to Yiotis *et al.* (2007) for the detailed derivation of Equation (6) and (7). Equation (7) suggests that $k_{r,w}$ is always bounded between 0 to 1 as it is only function of S_w . However, $k_{r,nw}$ may become greater than 1 in case $M > 1$ due to the 'lubricating' effect of the wetting fluid.

Simulations have been performed for $M = 0.01, 1$ and 100 with S_w ranging from 0 to 1 with increment of 0.1. Total 100 grid cells have been taken across L . Figure 2 shows analytical and numerical velocity profiles for the case of $M = 100$ and $S_w = 0.5$. Figure 3 compares analytical and numerical relative permeabilities for wetting and non-wetting fluids for $M = 100$. Our simulations show an excellent agreement with analytical results having maximum error in relative permeability to be less than 0.2%. Huang and Lu (2009) reports this error to be nearly 7% using multiphase Lattice-Boltzmann method. Also, in their results, continuity of the shear stress is not maintained at the interface and hence a velocity jump is observed.

Viscous Finger in a 2D Channel

A viscous finger is an instability that may occur when a low viscosity fluid displaces high viscosity fluid. In this section, we simulate the single viscous finger formation phenomenon in a 2D channel. Initially, a channel with finite width H is fully saturated with high viscosity fluid 2. Fluid 1 with low

viscosity is introduced in x -direction with fully developed velocity profile and displaces Fluid 2 from channel. During this displacement process the viscous finger gets developed and produces a steady state shape (constant finger tip velocity) as shown in Figure 4.

Average velocity of the inlet fluid 1 is U_1 . Viscosity ratio $M = \mu_2/\mu_1$ is 20 which is generally experienced during water flooding process where water displaces oil. No-slip boundary conditions are applied at the channel walls in y -direction whereas velocity inlet and pressure outlet boundary conditions are applied in x -direction. Densities of fluid 1 and 2 have been chosen equal ($\rho_1 = \rho_2 = \rho$) for the simplicity as the density ratio does not affect the finger formation and

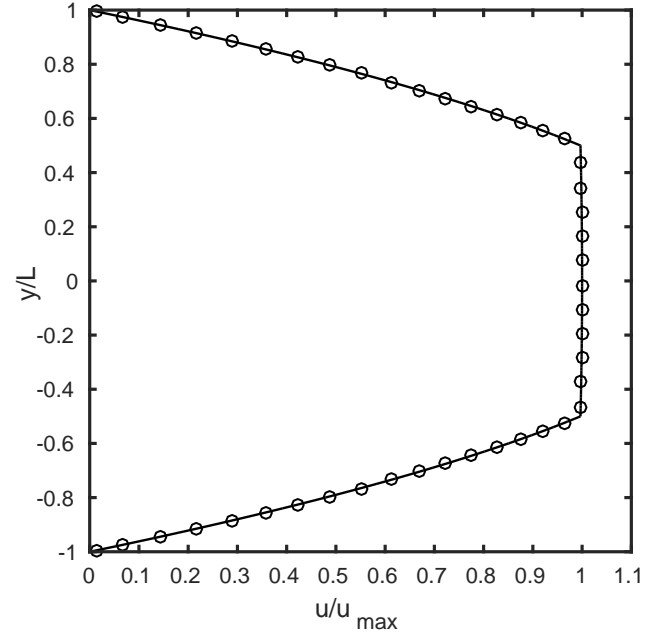


Figure 2: Analytical (—) and numerical (○) velocity profiles for $M = 100$ and $S_w = 0.5$.

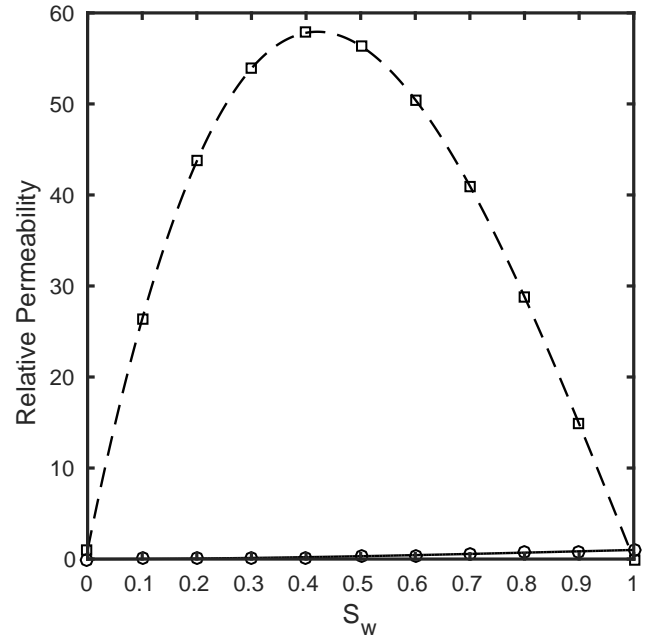


Figure 3: Relative permeability with S_w for $M = 100$: analytical $k_{r,w}$ (—), numerical $k_{r,w}$ (○), analytical $k_{r,nw}$ (---) and numerical $k_{r,nw}$ (□).

its properties. Capillary and Reynolds number defined using the steady state finger tip velocity U_t are $Ca = \mu_2 U_t / \sigma$ and $Re = \rho U_t H / \mu_2$. From the mass balance, one can find $U_t = U_1 H / W$. To maintain the flow in Stokes regime, value of Ca and Re have been chosen such that $ReCa = 10^{-3}$.

Simulations have been performed with different Ca ranging from 0.025 to 3 with 64 grid cells across the height H . The steady state finger width W and finger tip radius R have been obtained. Figure 5 compares the same finger characteristics parameters in dimensionless form with the results obtained using boundary element method by Halpern and Gaver (1994). Our results of W/H show an excellent match with maximum deviation to be less than 2%. Also, R/H shows an excellent match at low Ca . However at higher Ca , deviation is higher due to higher tip curvatures.

Contact Force Calculation

In this section, we present the grid independence study for contact force at three phase contact line in 3D. Initially, an oil droplet of equivalent radius $R_{eq} = 1$ mm (Volume $V = \frac{4}{3}\pi R_{eq}^3$) is placed on the solid sphere of radius $R_s = 1$ mm such that it inscribes $\theta = 60^\circ$ as shown in Figure 6. In this position,

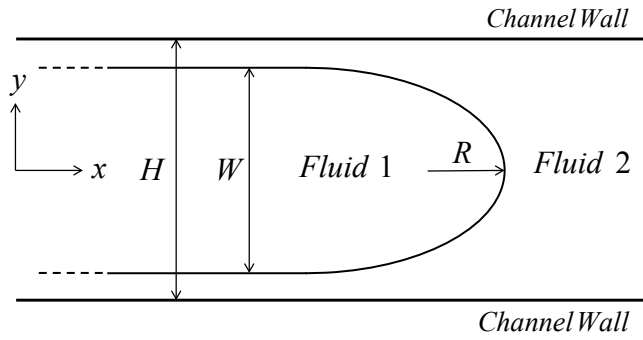


Figure 4: Schematic diagram of steady state viscous finger in 2D channel.

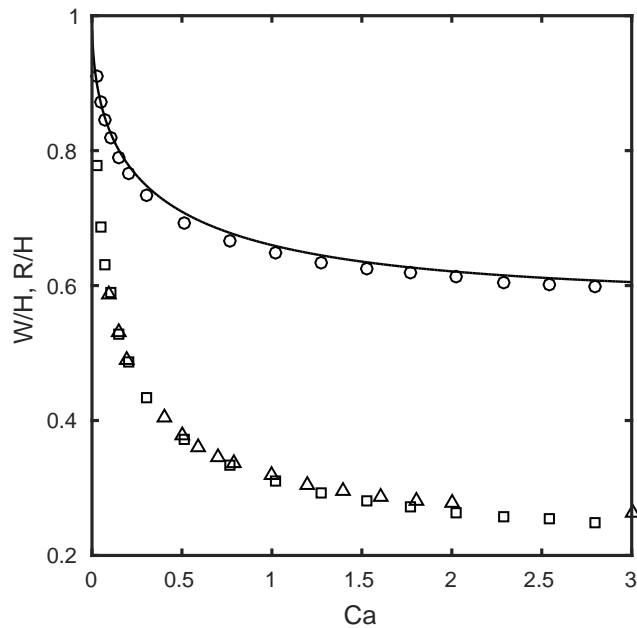


Figure 5: Dimensionless characteristics parameters of steady state viscous finger in 2D channel with Ca a) finger width W/H : present (\circ), Halpern and Gaver (1994) ($-$) b) finger tip radius R/H : present (\square), Halpern and Gaver (1994) (\triangle).

droplet radius $R_d = 1.1082R_{eq}$, contact radius $R_c = 0.9069R_{eq}$ and $\psi = 35.08^\circ$ (Patel *et al.*, 2017). The oil droplet is surrounded by water with $\sigma = 0.03$ N/m. Using a force balance, contact force can be given by following expression,

$$\begin{aligned} F_{y,analytical} &= 2\pi\sigma R_c \sin\psi \\ F_{x,analytical} &= F_{z,analytical} = 0 \end{aligned} \quad (8)$$

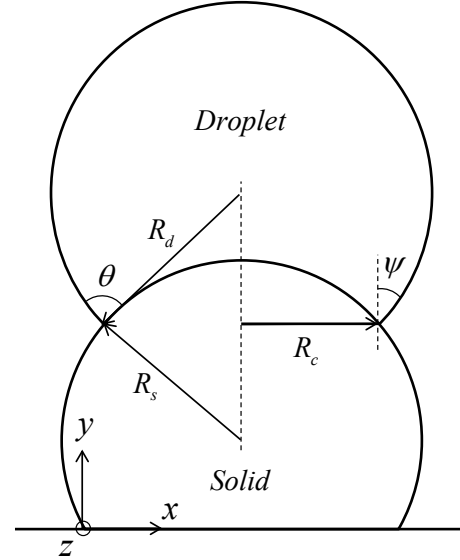


Figure 6: Schematic diagram of droplet on solid sphere for contact force calculation (xy cross section).

Numerical values of the contact force in all three directions have been calculated using the method proposed by Washino *et al.* (2013). The relative error in the contact force in the y -direction is given by following expression,

$$Error(\%) = \frac{|F_{y,analytical} - F_{y,numerical}|}{F_{y,analytical}} \times 100\% \quad (9)$$

Figure 7 shows the relative error in the contact force in y -direction with different grid resolutions. It follows a 1^{st} order trend. However, even at lower grid resolution ($R_c/\Delta \approx 5$) the relative error is lower than 4%. Moreover, the maximum value of the numerical contact forces in x - and z -directions (F_x/F_y and F_z/F_y) is less than 10^{-6} which affirms the accuracy of present coupled IBM-VOF implementation.

RESULTS

In this section, we present pore-scale simulations of a water flooding process such as encountered in enhanced oil recovery. Our aim is to investigate the wettability effects on the mobility of oil through oil-wet and neutrally-wet rocks.

The physical properties of an oil-water system are: $\mu_{oil} = 0.02$ Pa·s, $\mu_{water} = 0.001$ Pa·s, $\rho_{water} = \rho_{oil} = 1000$ kg/m³ and $\sigma = 0.03$ N/m. Initial distribution of oil and water phases in homogeneous configuration of spheres (rock structure) is shown in Figure 8a. Saturation of oil $S_{oil} = 0.5$ and structure porosity $\phi = 0.5$. Wettability is altered by varying θ for oil-wet ($\theta < 90^\circ$) and neutrally-wet ($\theta = 90^\circ$) rocks.

Simulations are performed on a 3D periodic domain with $100 \times 100 \times 100$ grid cells of size $\Delta = 1.5 \times 10^{-5}$ m. In this case, the number of grid cells across the radius of the sphere is around 39 which is quite sufficient to resolve accurate contact force. A constant body force of 10^5 N/m³ is applied in

x -direction to both oil and water phases. Superficial velocities of both phases are monitored to decide if the flow has reached steady state. Figure 9 shows the plot of superficial velocity of oil U_{oil} with time for oil-wet rocks with $\theta = 45^\circ$. One can observe that U_{oil} repeats itself after a fixed time interval and hence the flow has reached (pseudo) steady state. A snapshot of this (pseudo) steady state is presented in Figure 8b. Using the time averaged superficial velocity of oil \bar{U}_{oil} , one can calculate Re , Ca and $k_{r,oil}$ by means of the following expressions,

$$\begin{aligned} Re &= \frac{\rho_{oil} \bar{U}_{oil} D}{\mu_{oil}} \\ Ca &= \frac{\mu_{oil} \bar{U}_{oil}}{\sigma} \\ k_{r,oil} &= \frac{\bar{U}_{oil}(S_{oil})}{\bar{U}_{oil}(S_{oil} = 1)} \end{aligned} \quad (10)$$

where D is the diameter of sphere and $\bar{U}_{oil}(S_{oil})$ is \bar{U}_{oil} at given S_{oil} . Maximum value of Re and Ca numbers among all simulations are 0.5 and 0.006 respectively ensuring the flow in the Stokes regime.

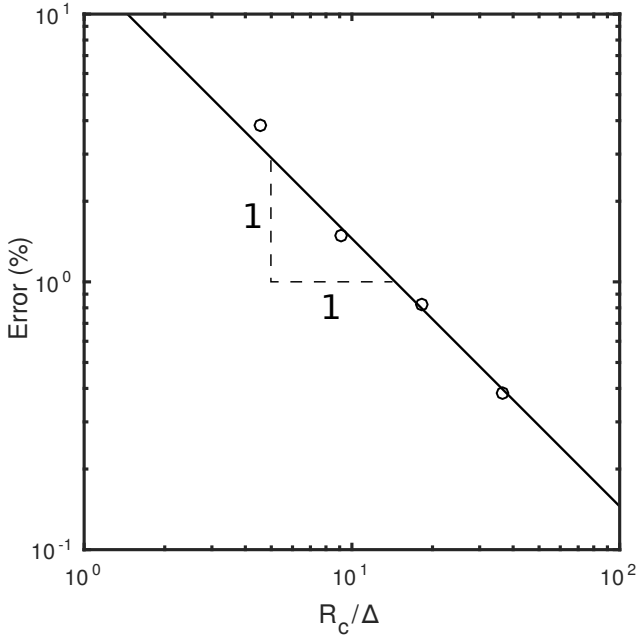


Figure 7: Grid independence study for contact force in y -direction: present (○) and 1st order line (–) for reference

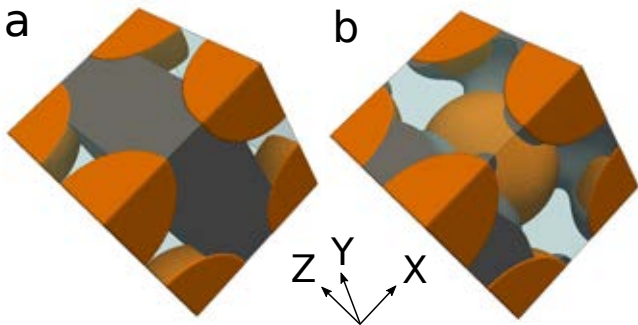


Figure 8: (a) Initial and (b) (pseudo) steady state fluid phase distribution for oil-wet rocks with $\theta = 45^\circ$, $\phi = 0.5$ and $S_{oil} = 0.5$. Oil, water and rocks are represented by grey, transparent blue and brown color respectively.

Figure 10 shows the $k_{r,oil}$ with θ for oil-wet ($\theta < 90^\circ$) and neutrally-wet ($\theta = 90^\circ$) rocks. Oil is more mobile in the neutrally-wet rocks as oil and water both have equal wettability towards rocks. As the rocks become more and more oil-wet ($\theta \rightarrow 0^\circ$), the tendency of oil to adhere with rocks increases and hence its mobility decreases. Interestingly, this decrement is non-linear and it may fully choke the flow at a lower body force. The maximum difference in $k_{r,oil}$ compared to its mean value is around 50% which suggests that the wettability has a strong effect on the oil recovery during water flooding process.

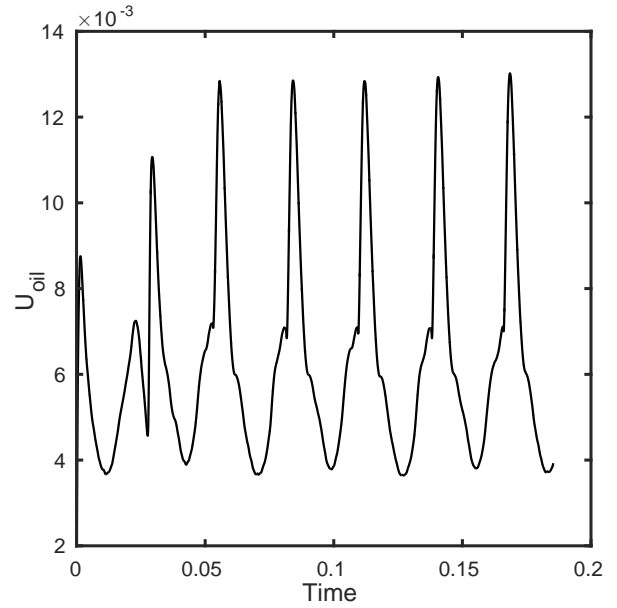


Figure 9: U_{oil} with time for oil-wet rocks with $\theta = 45^\circ$, $\phi = 0.5$ and $S_{oil} = 0.5$.

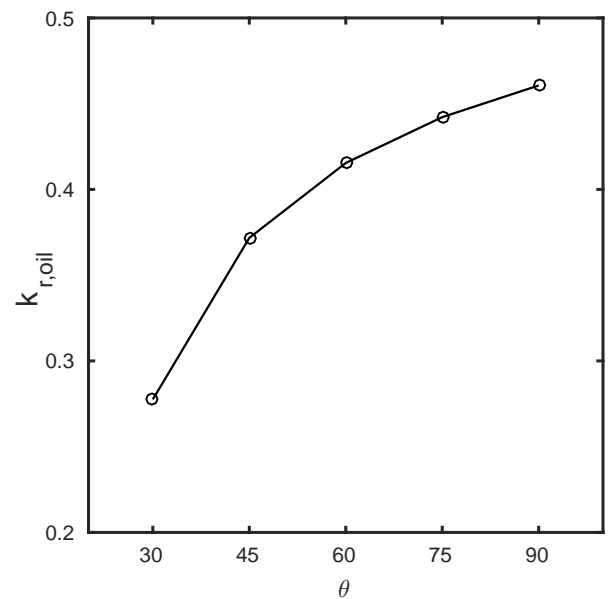


Figure 10: $k_{r,oil}$ with θ for oil-wet ($\theta < 90^\circ$) and neutrally-wet ($\theta = 90^\circ$) rocks with $\phi = 0.5$ and $S_{oil} = 0.5$

CONCLUSION

In this paper a methodology is presented to simulate the multiphase flows involving oil and water from complex geometries. IBM is used to simulate the dynamic interactions between fluids and complex geometries on a Cartesian computational grid. Present IBM is second order, direct forcing, implicit and doesn't require any calibration for different geometries. Sharp interface VOF method is used to track multi-fluid interfaces. IBM and VOF are coupled at the three phase contact line via the apparent contact angle. Present methodology works perfectly well for the validation/verification test cases involving oil-water flows and yields more accurate results compared to simulations performed using the Lattice-Boltzmann method. Pore-scale water flooding simulations are presented to quantify the effect of wettability on the mobility of oil through oil-wet and neutrally-wet homogeneous rocks. Obtained results show that the oil-wet rocks can drastically reduce the mobility of oil.

In future, this work will be continued on large randomized structures with number of spheres in order of 100 to quantify the effect of wettability, porosity, capillary number, saturation, viscosity ratio etc. on the mobility of oil-water flows.

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